

*Springs :*

$$y = A \cos \omega t;$$

$$\omega^2 = \frac{k}{m};$$

$$E = \frac{1}{2} kx^2 + \frac{1}{2} mv^2$$

$$\omega = 2\pi f = \frac{2\pi}{T}$$

*Simple pendulum :*

(0.1)

$$\omega^2 = \frac{g}{l}$$

*Waves :*

$$y = A \sin(kx - \omega t); k = \frac{2\pi}{\lambda}; \omega = \frac{2\pi}{T}; v = \frac{\lambda}{T} = \lambda f = \frac{\omega}{k}$$

(0.2)

$$\text{transversal waves on strings : } v = \sqrt{\frac{F_T}{m/L}} = \sqrt{\frac{F_T}{\mu}}$$

$$\text{sound-waves in water and other mediums with density } \rho : v = \sqrt{\frac{B}{\rho}}$$

Power and intensity of waves:

*Waves on strings :*

$$\text{Power} = \frac{1}{2} \mu v (\omega y_{\max})^2$$

*Sound waves :*

$$I = \frac{1}{2} \rho v (\omega s_{\max})^2; \rho = \text{density of the medium (air, water, concrete, steel etc.)}$$

Superposition (addition) of trigonometric functions:

$$\text{soundlevel } \beta = 10 \log \frac{I}{I_0}$$

Interference of sine waves with the same amplitude but different arguments:

$$A_1 \sin \theta_1 + A_1 \sin \theta_2 = 2A_1 \cos \frac{\theta_2 - \theta_1}{2} \sin \frac{\theta_2 + \theta_1}{2}$$

Double slit experiment:

If two waves differ only by a constant phase-shift they show the effect of destructive and constructive interference:

$$y_1 + y_2 = 2A \cos \frac{\phi}{2} \cdot \sin \left( kx - \omega t + \frac{\phi}{2} \right)$$

(0.3) In the DSE the phaseshift is produced by different path lengths:

$$y_2 = A \sin [k(x + \Delta x) - \omega t]; \Delta x = d \sin \theta$$

Constructive interference occurs when  $\Delta x = n\lambda$

### Standing Waves:

**When waves occupy the same space superposition can form an interference pattern. Under certain conditions they form standing waves. Standing waves on a string under tension, for example occur on string instruments guitars, violins, cellos, etc. Two waves traveling in opposite directions interfere according to:**

$$A \sin(kx - \omega t) + A \sin(kx + \omega t) = 2A \cos(\omega t) \sin kx$$

(0.4) If the string has length L and is tied down at both ends, we must have a 0 value for  $x=L$

$$\sin kL = 0 \Rightarrow kL = n\pi \Rightarrow \lambda_n = \frac{2L}{n}; f_n = \frac{v}{2L} n; n \text{ even and odd integers.}$$

If one end is closed and the other end open, the open end at  $x=L$  requires that:

$$\sin(kL) = \frac{2n+1}{2} \pi; n = 0, 1, 2, 3$$

$$\frac{2\pi}{\lambda_{2n+1}} L = \frac{2n+1}{2} \pi \Rightarrow \lambda_{2n+1} = \frac{4L}{2n+1}$$

(0.5)  $f_{2n+1} = \frac{v}{4L} (2n+1) \Rightarrow f_1 = \frac{v}{4L}; f_3 = 3f_1; f_5 = 5f_1$  the even overtones are missing.

$$\text{example: } v = 343 \frac{m}{s}; L = 0.2 \text{ meters} \Rightarrow f_1 = \frac{343 \frac{m}{s}}{0.2m} = 1715 \text{ Hz}; \lambda_1 = 0.8m; \lambda_3 = \frac{0.8}{3} m$$

## Doppler effect: source moves, receiver is stationary;

Doppler effect for sound, source moves, observer is stationary:

$$a) \lambda' = \lambda \left( 1 - \frac{s}{v} \right) \text{ source approaching with speed } s$$

$$(0.6) b) \lambda' = \lambda \left( 1 + \frac{s}{v} \right) \text{ source receding with speed } s$$

To get from wavelengths  $\lambda$  to frequencies  $f$ , we just remember the general relationship  $\lambda f = v$  and  $\lambda' f' = v$ ; where  $v$  is the speed of the wave.

## Doppler effect: receiver moves, source is stationary

The relationship between speed of sound and wavelength is in the reference frame of the source:

$$v = \lambda f \text{ or } a) f = \frac{v}{\lambda}$$

In the reference frame of the moving receiver, it is

$$b) v' = \lambda f' \text{ or } a) f' = \frac{v'}{\lambda} \Rightarrow \frac{f'}{f} = \frac{v'}{v}$$

As we have  $v' = v \pm V_R$ , we get:

$$c) f' = f \left( \frac{v \pm V_R}{v} \right) = f \left( 1 \pm \frac{V_R}{v} \right)$$

The frequency increases when the receiver moves towards the stationary source.

## Relativistic Doppler effect for light:

(0.7) Doppler effect for light (relativistic):

$$a) \lambda' = \lambda \sqrt{\frac{1 - \frac{v}{c}}{1 + \frac{v}{c}}}; f' = f \sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}} \qquad b) \lambda' = \lambda \sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}}; f' = f \sqrt{\frac{1 - \frac{v}{c}}{1 + \frac{v}{c}}}$$

a) for a star approaching with the speed  $v$ ; blue shift

b) for a star receding with the speed  $v$ ; red shift

**Hubble** constant: Once we know the speed of a galaxy we can approximate the distance  $R$  of the galaxy to our own with **Hubble's** formula:

$$\begin{aligned} R &= \frac{v}{H} \\ (0.8) \quad H &= \frac{0.017m}{s \cdot \text{lightyears}} \end{aligned}$$

$$\begin{aligned} R &= \frac{v}{H} \\ (0.9) \quad H &= \frac{0.017m}{s \cdot \text{lightyears}} \end{aligned}$$